

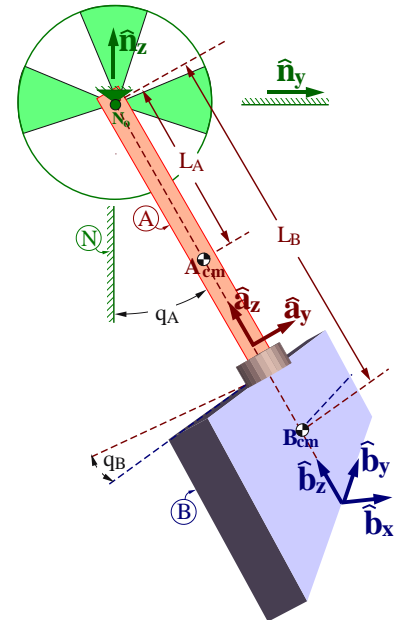
# Motivating example (MIPSI): Babyboot

## Model

The figure to the right is a schematic representation of a swinging babyboot attached by a shoelace to a rigid support. The mechanical model of the babyboot consists of a thin uniform rod  $A$  attached to a fixed support  $N$  by a revolute joint, and a uniform plate  $B$  connected to  $A$  with a second revolute joint so that  $B$  can rotate freely about  $A$ 's axis.

Note: The revolute joints' axes are *perpendicular*, not parallel.

- **Bodies:** The rod and plate are **rigid** (inflexible/undeformable).
- **Connections:** The revolute joints are **ideal** (massless, frictionless, with no slop or flexibility).
- **Force:** **Earth's gravity** is uniform and constant. Other contact forces (e.g., **air resistance, solar/light pressure**) and distance forces (e.g., **electromagnetic, other gravitational**) are negligible.
- **Newtonian reference frame:** **Earth**



## Identifiers

Right-handed sets of unit vectors  $\hat{n}_x, \hat{n}_y, \hat{n}_z$ ;  $\hat{a}_x, \hat{a}_y, \hat{a}_z$ ; and  $\hat{b}_x, \hat{b}_y, \hat{b}_z$  are fixed in  $N$ ,  $A$ , and  $B$ , respectively, with  $\hat{n}_x = \hat{a}_x$  parallel to the revolute axis joining  $A$  to  $N$ ,  $\hat{n}_z$  vertically upward,  $\hat{a}_z = \hat{b}_z$  parallel to the rod's long axis (and the revolute axis joining  $B$  to  $A$ ), and  $\hat{b}_z$  perpendicular to plate  $B$ .

Quantity	Symbol	Type	Value
Earth's gravitational constant	$g$	Constant	9.81 m/s <sup>2</sup>
Distance between $N_o$ and $A_{cm}$	$L_A$	Constant	7.5 cm
Distance between $N_o$ and $B_{cm}$	$L_B$	Constant	20 cm
Mass of $A$	$m^A$	Constant	0.01 kg
Mass of $B$	$m^B$	Constant	0.1 kg
$A$ 's moment of inertia about $A_{cm}$ for $\hat{a}_x$	$I^A$	Constant	0.05 kg*cm <sup>2</sup>
$B$ 's moment of inertia about $B_{cm}$ for $\hat{b}_x$	$I_x^B$	Constant	2.5 kg*cm <sup>2</sup>
$B$ 's moment of inertia about $B_{cm}$ for $\hat{b}_y$	$I_y^B$	Constant	0.5 kg*cm <sup>2</sup>
$B$ 's moment of inertia about $B_{cm}$ for $\hat{b}_z$	$I_z^B$	Constant	2.0 kg*cm <sup>2</sup>
Angle from $\hat{n}_z$ to $\hat{a}_z$ with $+\hat{n}_x$ sense	$q_A$	Dependent variable	Varies
Angle from $\hat{a}_y$ to $\hat{b}_y$ with $+\hat{a}_z$ sense	$q_B$	Dependent variable	Varies
Time	$t$	Independent variable	Varies

## Physics

Physics from [www.MotionGenesis.com](http://www.MotionGenesis.com) ⇒ [Get Started](#) ⇒ Chaotic Pendulum (Babyboot).

The ODEs (ordinary differential equations) governing the motion of this mechanical system are<sup>11</sup>

$$\ddot{q}_A = \frac{2 \dot{q}_A \dot{q}_B \sin(q_B) \cos(q_B) (I_x^B - I_y^B) - (m^A L_A + m^B L_B) g \sin(q_A)}{I^A + m^A L_A^2 + m^B L_B^2 + I_x^B \cos^2(q_B) + I_y^B \sin^2(q_B)}$$

$$\ddot{q}_B = \frac{-\dot{q}_A^2 \sin(q_B) \cos(q_B) (I_x^B - I_y^B)}{I_z^B}$$

## Simplify and solve

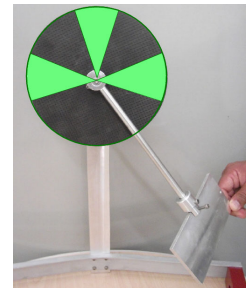
The set of differential equations governing the babyboot's motion are (circle the appropriate qualifiers)

Uncoupled	Linear	Homogeneous	Constant-coefficient	1 <sup>st</sup> -order
Coupled	Nonlinear	Inhomogeneous	Variable-coefficient	2 <sup>nd</sup> -order

Computers has revolutionized the solution of differential equations. There are a many numerical algorithms for solving nonlinear, coupled, variable coefficient, ODEs (ordinary differential equations) including Euler's method, predictor-corrector, Runge-Kutta, etc. In addition, there are many programs (MATLAB<sup>®</sup>, MotionGenesis, WolframAlpha, etc.) that make it easy to solve ODEs.

### Computer (numerical) solution of ODEs with MotionGenesis (with plotting)

```
Variable qA'', qB''      % Angles and first/second time-derivatives.
%-----
qA'' = 2*( 508.89*sin(qA) - sin(qB)*cos(qB)*qA'*qB' ) / (-21.556 + sin(qB)^2)
qB'' = -sin(qB)*cos(qB)*qA'^2
%-----
Input  tFinal = 10 sec, tStep = 0.02 sec, absError = 1.0E-07
Input  qA = 90 deg, qB = 1.0 deg, qA' = 0.0 rad/sec, qB' = 0.0 rad/sec
OutputPlot t sec, qA degrees, qB degrees
%-----
ODE() solveBabybootODE
Quit
```



### Alternately: Simplify via linearization and solve analytically (valid only for very small angles)

Linearizing these ODEs about  $q_A = 0$  and  $q_B = 0$  produces a simpler set of ODEs, namely

$$\ddot{q}_A = -\omega^2 q_A \quad \ddot{q}_B = 0 \quad \text{where } \omega = \sqrt{\frac{(m^A L_A + m^B L_B) g}{I^A + m^A L_A^2 + m^B L_B^2 + I_x^B}}$$

When released from rest [i.e., no initial spin, i.e.,  $\dot{q}_A(0) = \dot{q}_B(0) = 0$ ], the solutions to these ODEs are

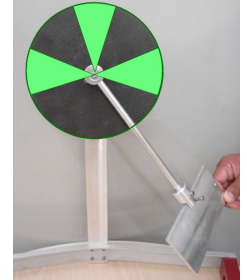
$$q_A(t) = q_A(0) \cos(\omega t) \quad q_B(t) = q_B(0) \quad \text{constant!}$$

<sup>11</sup>Four methods for forming equations of motion are: **Free-body diagrams** of  $A$  and  $B$  (which is inefficient as it introduces up to 10 unknown force/torque measures); D'Alembert's method (**road maps** of Section 22.6) which efficiently forms the two equations shown for  $\ddot{q}_A$  and  $\ddot{q}_B$  (but require a clever selection of systems, points, and unit vectors); **Lagrange's equations** (an energy-based method that automates D'Alembert's cleverness); **Kane's equations** (a modern efficient blend of D'Alembert and Lagrange).

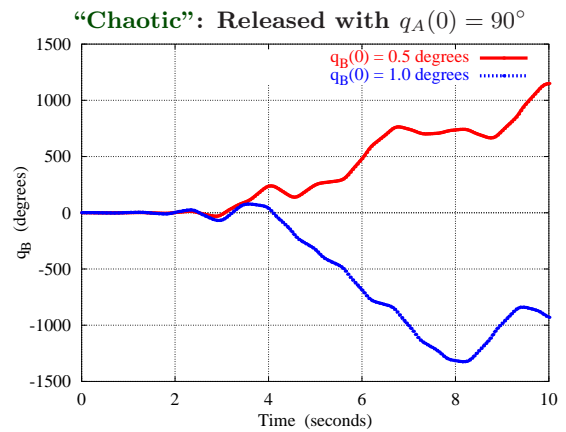
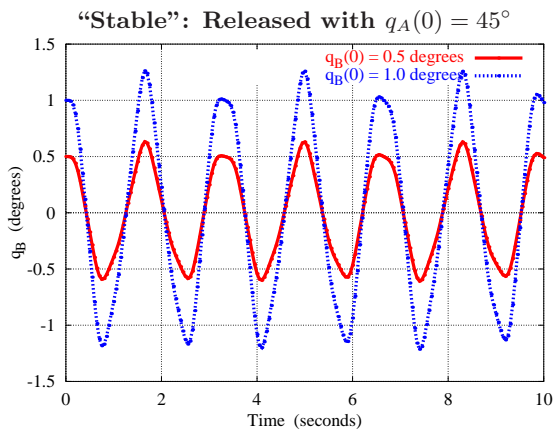
## Interpret

The solution to these differential equations reveals this simple system has strange, non-intuitive motion.<sup>12</sup> For certain initial values of  $q_A$ , the motion of plate  $B$  is well-behaved and “stable”. Alternately, for other initial values of  $q_A$ ,  $B$ ’s motion is “**chaotic**” – meaning that a small variation in the initial value of  $q_B$  or numerical integration inaccuracies lead to dramatically different results (these ODEs are used to test the accuracy of numerical integrators – the plots below required a numerical integrator error of  $\text{absError} = 1 \times 10^{-7}$ ).

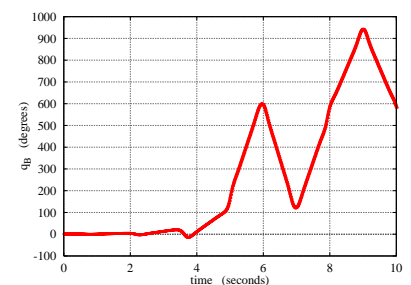
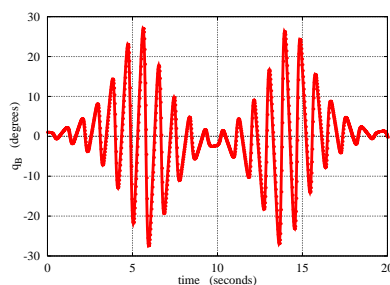
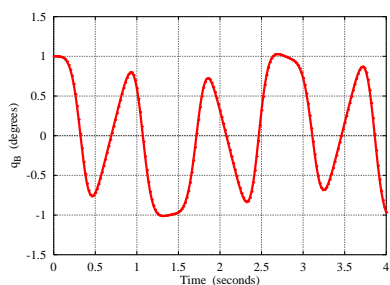
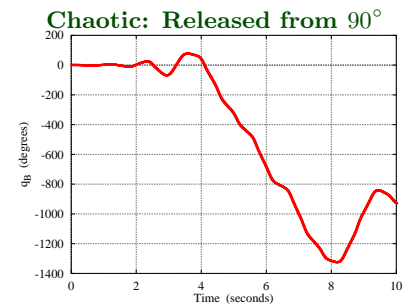
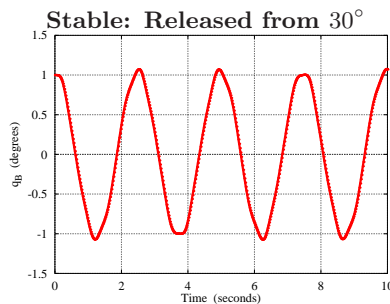
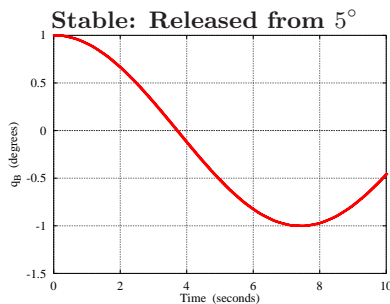
The following chart and figure to the right shows this system’s regions of stability (**black**) and instability (**green**). Notice the “**chaotic**” plot below shows  $q_B$  is *very* sensitive to initial values. A  $0.5^\circ$  change in the initial value of  $q_B(0)$  results in more than a  $2000^\circ$  difference in the value of  $q_B(t = 10)$ !



Initial value of $q_A$		Stability	
$0^\circ \leq q_A(0) \leq 71.3^\circ$		<b>Stable</b>	<b>black</b>
$71.4^\circ \leq q_A(0) \leq 111.77^\circ$		<b>Unstable</b>	<b>green</b>
$111.78^\circ \leq q_A(0) \leq 159.9^\circ$		<b>Stable</b>	<b>black</b>
$160.0^\circ \leq q_A(0) \leq 180.0^\circ$		<b>Unstable</b>	<b>green</b>



## Investigation of stability: More simulation results



<sup>12</sup>More information about this problem is in “Mechanical Demonstration of Mathematical Stability and Instability”, *International Journal of Engineering Education (Journal of Mechanical Engineering Education)*, Vol. 2, No. 4, 1974, pp. 45-47, by Thomas R. Kane. Or visit [www.MotionGenesis.com](http://www.MotionGenesis.com) ⇒ [Get Started](#) ⇒ [Chaotic Pendulum \(Babyboot\)](#).